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BIAXIAL PHASE TRANSITIONS IN NEMATIC LIQUID CRYSTALS



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Summary

*General quadrupolar interaction: Straley's potential
Mean-field model and Phase diagram for $\gamma = 0$*

Reference model parameters domain

Conjugated phase diagram

Extension to a mildly repulsive model

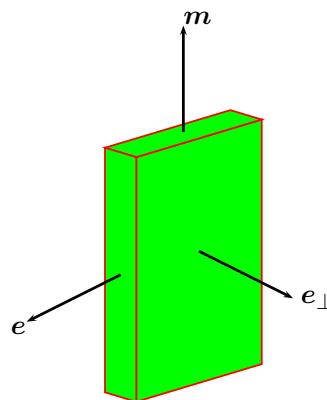
Molecular Biaxiality

We consider a homogeneous nematic liquid crystal intrinsically biaxial in the absence of any external field. It is made up of biaxial molecules which we can imagine as being described by a biaxial tensor decomposed into two traceless, irreducible orthogonal components

$$\mathbf{q} := \mathbf{m} \otimes \mathbf{m} - \frac{1}{3} \mathbf{I}$$

$$\mathbf{b} := \mathbf{e} \otimes \mathbf{e} - \mathbf{e}_\perp \otimes \mathbf{e}_\perp$$

\mathbf{m} , \mathbf{e} , \mathbf{e}_\perp axes of any molecular polarizability



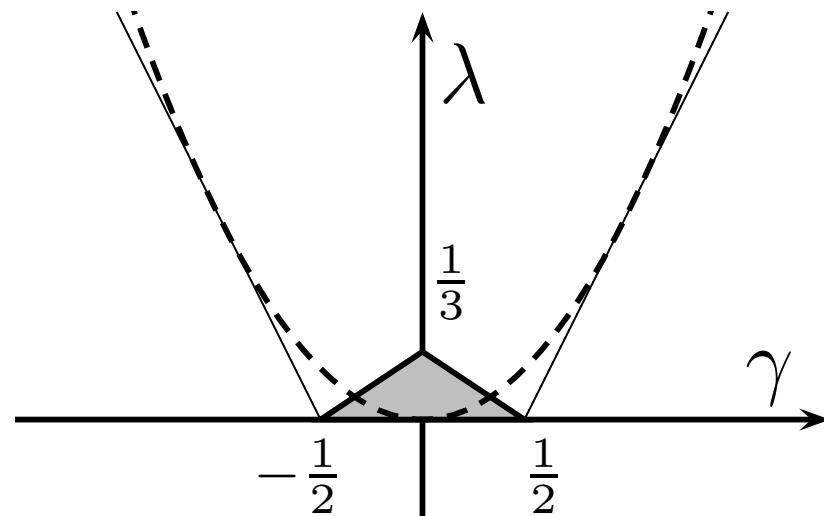
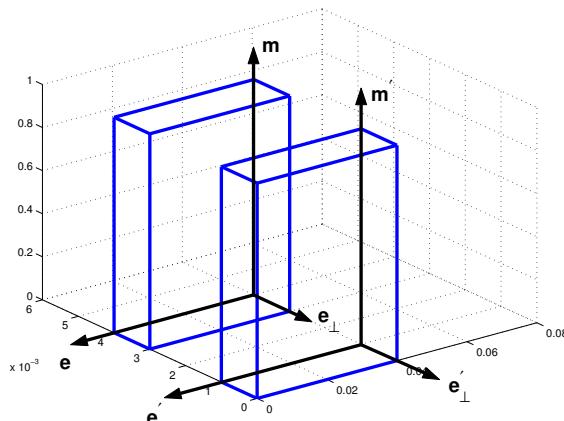
Interaction potential

Letting (\mathbf{q}, \mathbf{b}) and $(\mathbf{q}', \mathbf{b}')$ represent two interacting molecules, the most general quadrupolar interaction potential coupling them is

$$V = -U_0 \{ \mathbf{q} \cdot \mathbf{q}' + \gamma(\mathbf{q} \cdot \mathbf{b}' + \mathbf{b} \cdot \mathbf{q}') + \lambda \mathbf{b} \cdot \mathbf{b}' \}$$

$U_0 > 0$ and λ, γ model parameters

Ground State



Model $\gamma = 0$

$$V = -U_0 \{ \mathbf{q} \cdot \mathbf{q}' + \lambda \mathbf{b} \cdot \mathbf{b}' \}$$

$$\frac{V}{U_0} = - \left\{ - \left(\lambda + \frac{1}{3} \right) + (1 - \lambda) (\mathbf{m} \cdot \mathbf{m}')^2 + 2\lambda \left[(\mathbf{e}'_\perp \cdot \mathbf{e}_\perp)^2 + (\mathbf{e}' \cdot \mathbf{e})^2 \right] \right\}$$

$$\mathbf{Q} := \langle \mathbf{q} \rangle = S(\mathbf{e}_z \otimes \mathbf{e}_z - \frac{1}{3}\mathbf{I}) + T(\mathbf{e}_x \otimes \mathbf{e}_x - \mathbf{e}_y \otimes \mathbf{e}_y)$$

$$\mathbf{B} := \langle \mathbf{b} \rangle = S'(\mathbf{e}_z \otimes \mathbf{e}_z - \frac{1}{3}\mathbf{I}) + T'(\mathbf{e}_x \otimes \mathbf{e}_x - \mathbf{e}_y \otimes \mathbf{e}_y)$$

pseudo-potential

$$U = -U_0(\mathbf{q} \cdot \mathbf{Q} + \lambda \mathbf{b} \cdot \mathbf{B})$$

Mean-field equations
 partition and distribution functions

$$Z(\mathbf{Q}, \mathbf{B}, \beta, \lambda) = \int_{\mathbb{T}} \exp(\beta(\mathbf{Q} \cdot \mathbf{q} + \lambda \mathbf{B} \cdot \mathbf{b}))$$

$$f = \frac{1}{Z} \exp[\beta(\mathbf{q} \cdot \mathbf{Q} + \lambda \mathbf{b} \cdot \mathbf{B})]$$

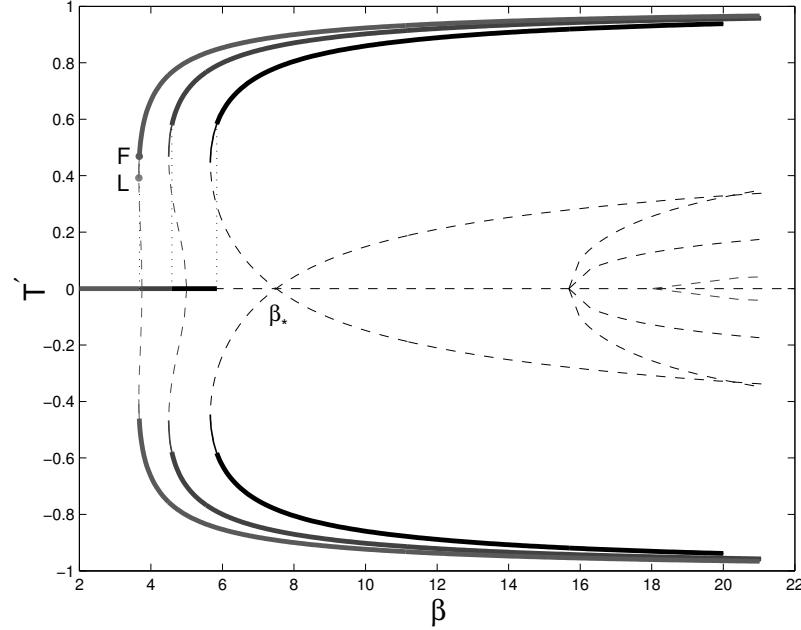
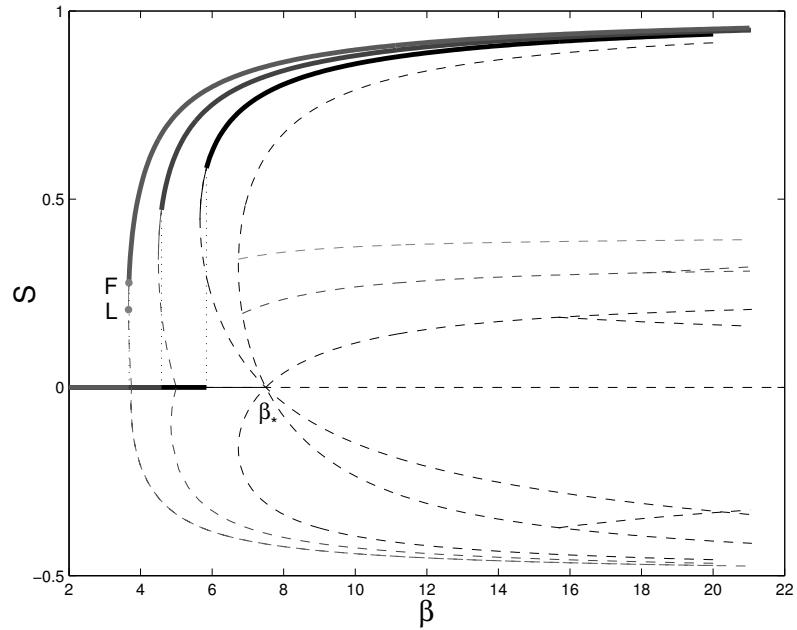
$$\mathbb{T} := \mathbb{S}^2 \times \mathbb{S}^1, \quad \beta := \frac{U_0}{k_B t}$$

k_B Boltzmann constant and t absolute temperature

$$\mathbf{Q} = \langle \mathbf{q} \rangle = \int_{\mathbb{T}} f \mathbf{q} \quad \mathbf{B} = \langle \mathbf{b} \rangle = \int_{\mathbb{T}} f \mathbf{b}$$

$$\mathcal{F}(\mathbf{Q}, \mathbf{B}, \beta, \lambda) = U_0 \left\{ \frac{1}{2} \mathbf{Q} \cdot \mathbf{Q} + \frac{\lambda}{2} \mathbf{B} \cdot \mathbf{B} - \frac{1}{\beta} \ln \left(\frac{Z(\mathbf{Q}, \mathbf{B}, \beta, \lambda)}{8\pi^2} \right) \right\},$$

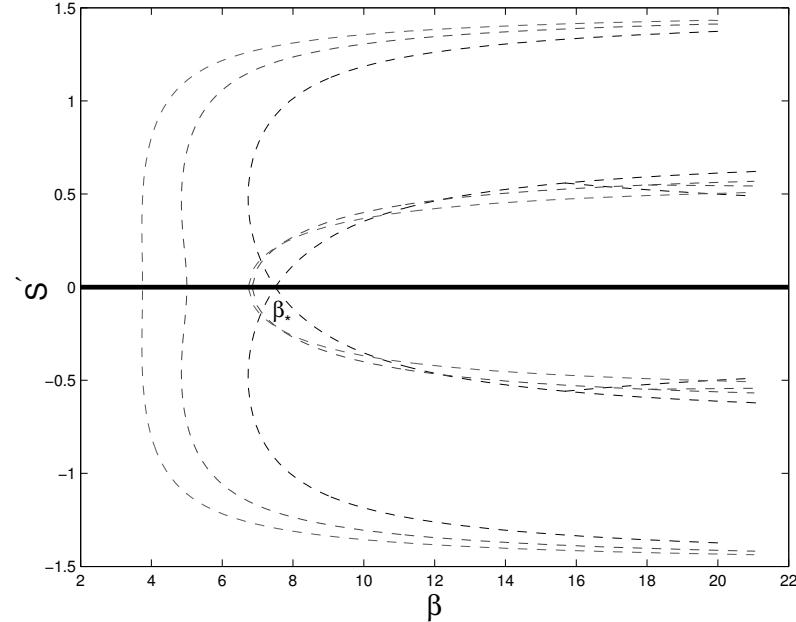
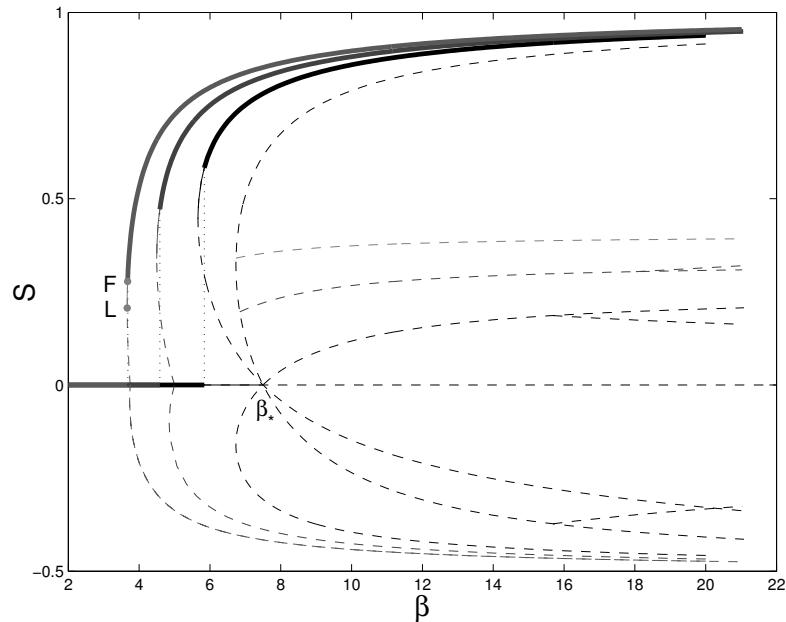
$$\lambda = 1/3, 1/2, 2/3$$



dominant stable biaxial states $(S, 0, 0, T')$

$$\mathbf{Q} = S(\mathbf{e}_z \otimes \mathbf{e}_z - \frac{1}{3}\mathbf{I}) \quad \mathbf{B} = T'(\mathbf{e}_x \otimes \mathbf{e}_x - \mathbf{e}_y \otimes \mathbf{e}_y)$$

$$\lambda = 1/3, 1/2, 2/3$$



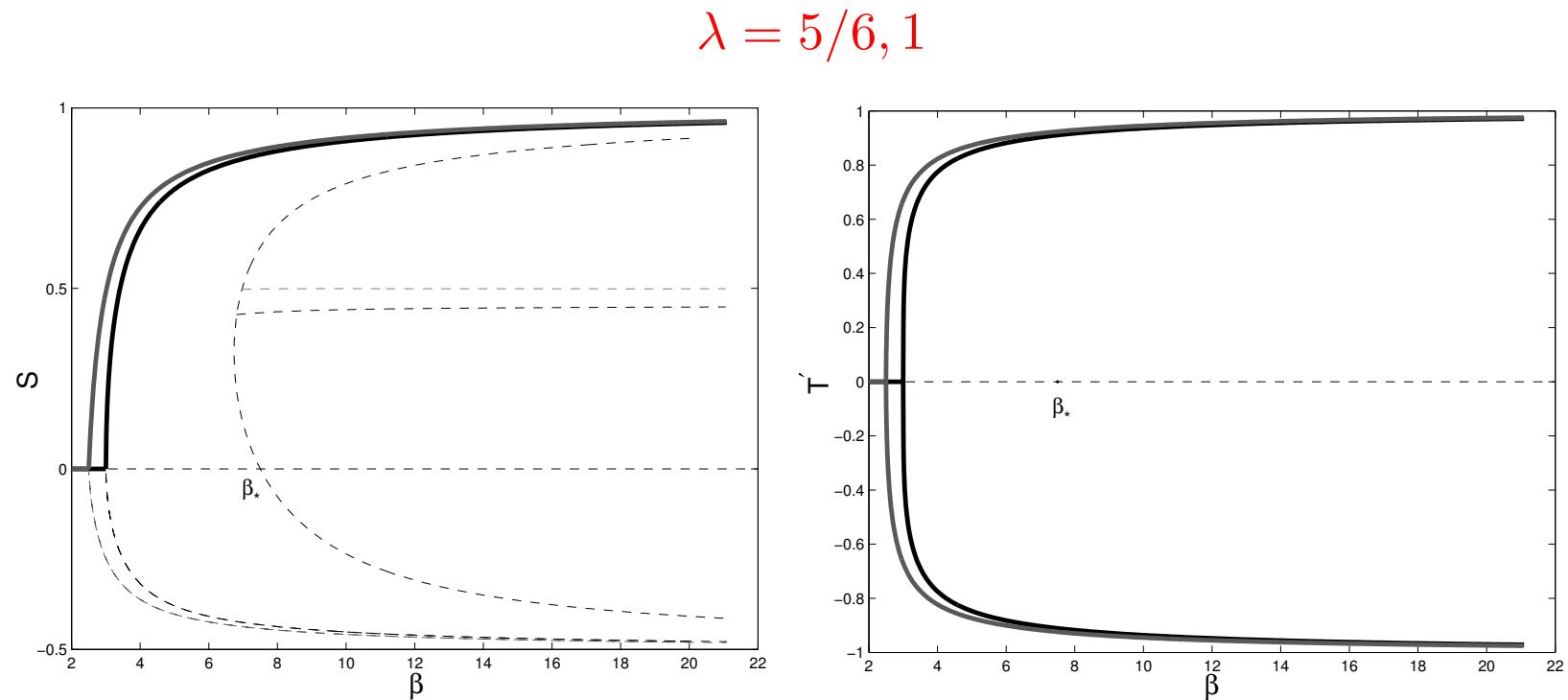
unstable uniaxial states with a greater value of free energy $(S, 0, S', 0)$

$$\mathbf{Q} = -S(\mathbf{e}_y \otimes \mathbf{e}_y - \frac{1}{3}\mathbf{I})$$

$$\mathbf{B} = -S'(\mathbf{e}_y \otimes \mathbf{e}_y - \frac{1}{3}\mathbf{I})$$

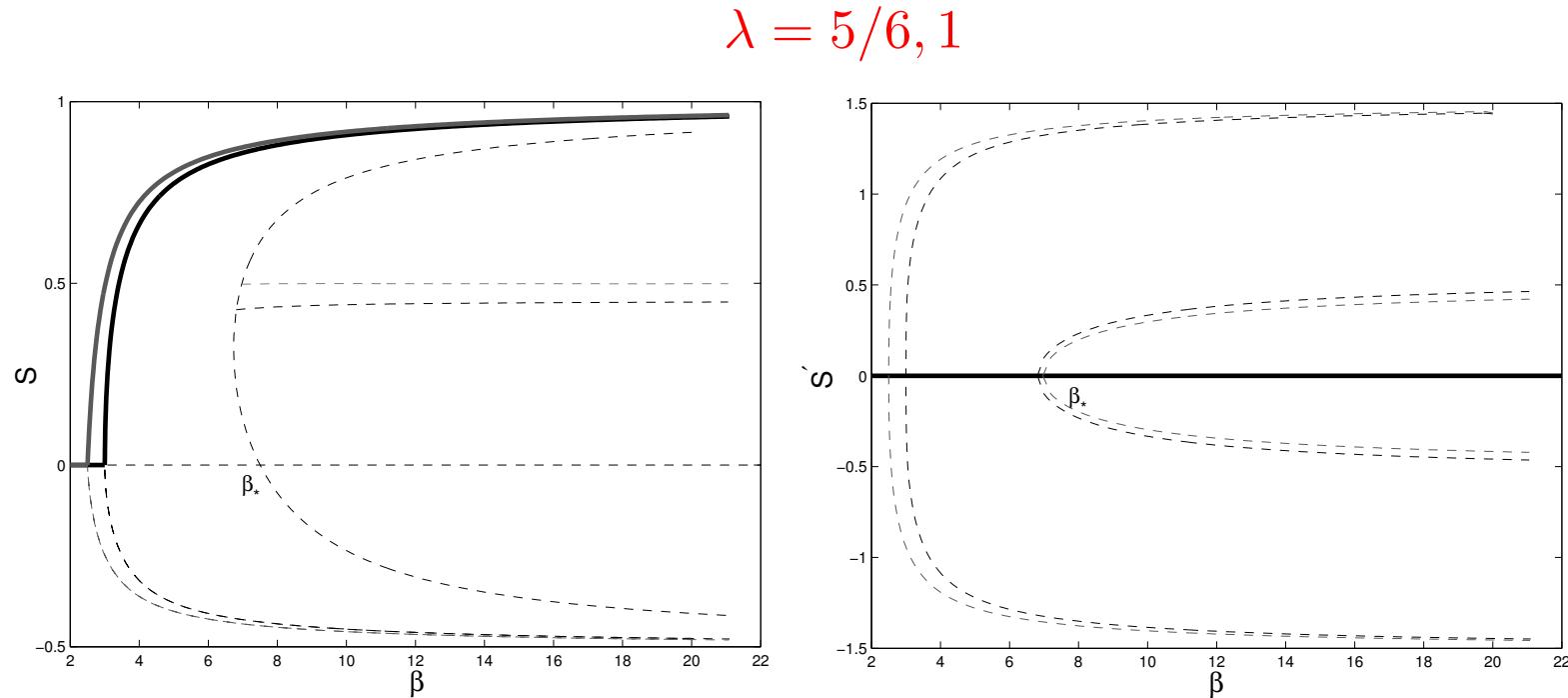
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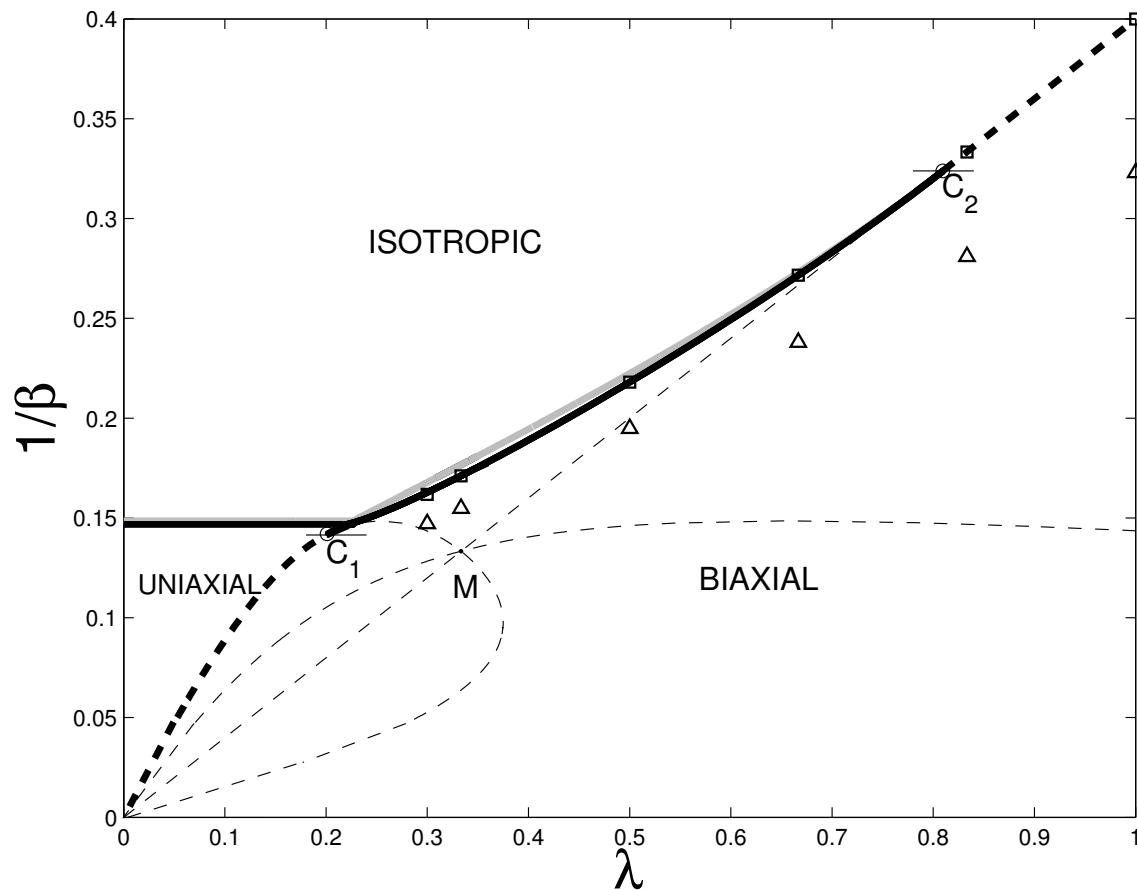
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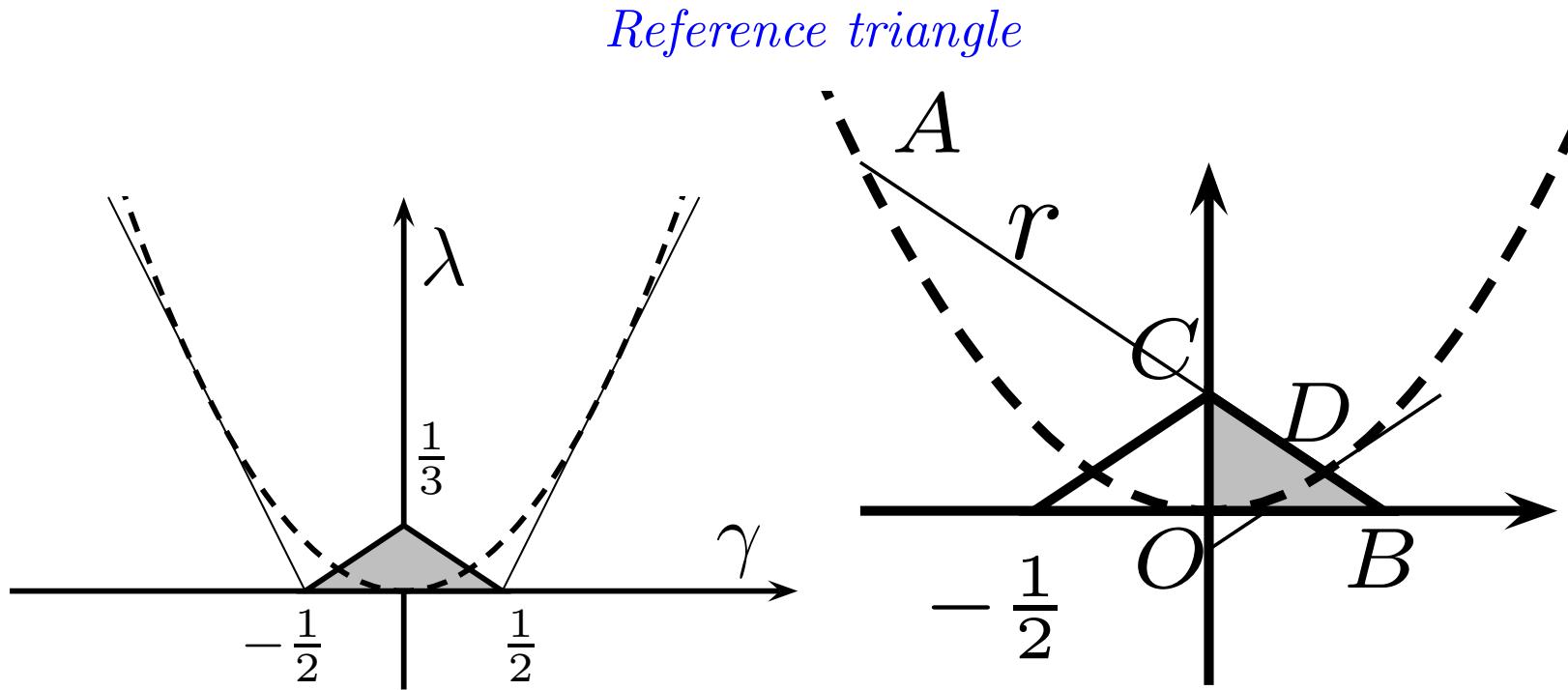
$$\mathbf{B} = -S'(\mathbf{e}_y \otimes \mathbf{e}_y - \frac{1}{3}\mathbf{I})$$

$$\mathbf{Q} = -S(\mathbf{e}_x \otimes \mathbf{e}_x - \frac{1}{3}\mathbf{I})$$

$$\mathbf{B} = -S'(\mathbf{e}_x \otimes \mathbf{e}_x - \frac{1}{3}\mathbf{I})$$

Phase Diagram and Tricritical points

Durand, Sonnet, Virga (2003) De Matteis, Romano ,Virga (2005)



$$D\left(\frac{1}{3}, \frac{1}{9}\right), \quad B\left(\frac{1}{2}, 0\right), \quad C\left(0, \frac{1}{3}\right), \quad A(-1, 1)$$

$$r : 2\gamma + 3\lambda - 1 = 0$$

Model: $2\gamma = 1 - 3\lambda$

$$\mathbf{q}^* := \mathbf{e} \otimes \mathbf{e} - \frac{1}{3}\mathbf{I}, \quad \mathbf{b}^* := \mathbf{m} \otimes \mathbf{m} - \mathbf{e}_\perp \otimes \mathbf{e}_\perp$$

$$V = -U_0 \left\{ \left(\frac{9\lambda - 1}{2} \right) \mathbf{q}^* \cdot \mathbf{q}^{*\prime} + \left(\frac{1 - \lambda}{2} \right) \mathbf{b}^* \cdot \mathbf{b}^{*\prime} \right\}$$

$$V^* = -U_0 \left(\frac{9\lambda - 1}{2} \right) \{ \mathbf{q}^* \cdot \mathbf{q}^{*\prime} + \lambda^* \mathbf{b}^* \cdot \mathbf{b}^{*\prime} \}$$

which is equivalent to $\gamma = 0$ model by transforming the model parameter λ and β as follows

$$\beta^* = \beta \left(\frac{9\lambda - 1}{2} \right), \quad \lambda^* = \frac{1 - \lambda}{9\lambda - 1}$$

in the (γ, λ) -plane

$$\left[\frac{1}{9}, 1 \right] \mapsto [0, \infty], \quad \left[\frac{1}{9}, \frac{1}{3} \right] \mapsto \left[\frac{1}{3}, \infty \right] \text{ i.e. } \overline{DA} \mapsto \overline{OP_\infty}, \overline{CD} \mapsto \overline{CP_\infty}$$

$$P_\infty (0, \infty)$$

Order parameters

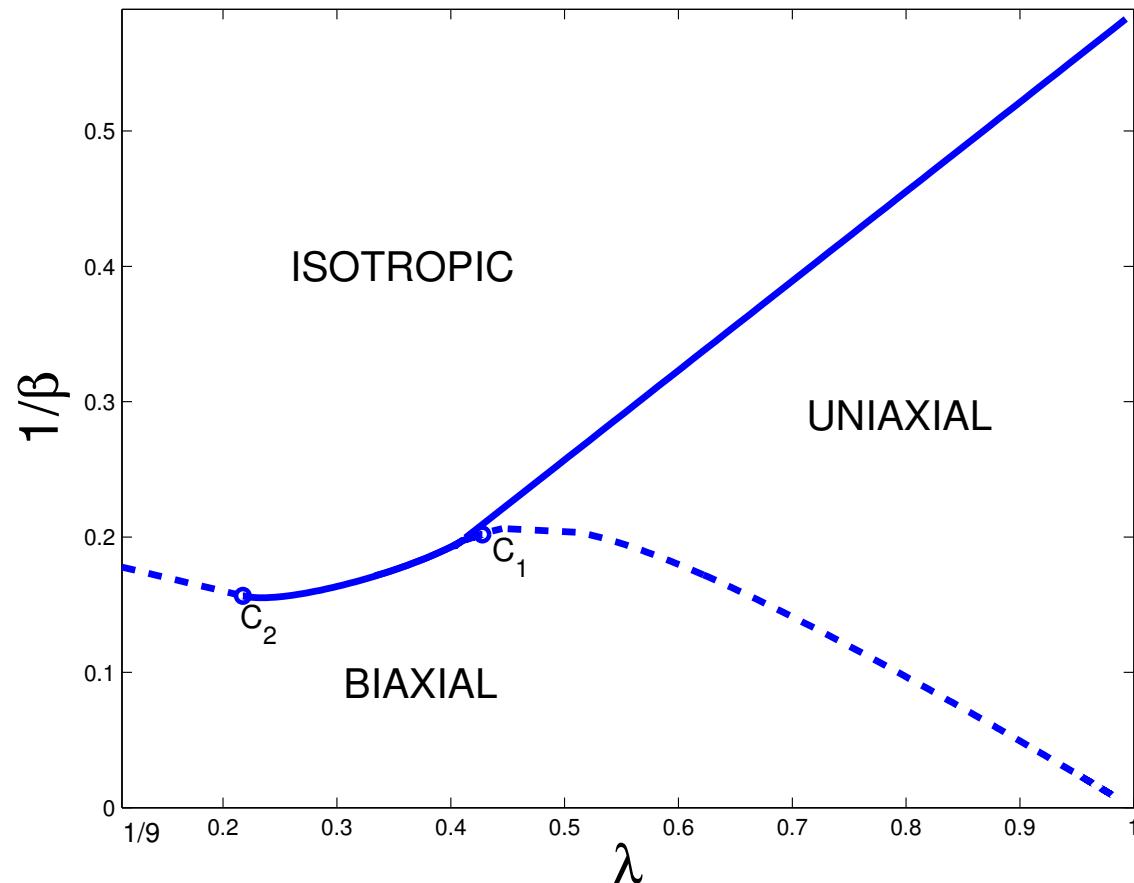
$$\mathbf{Q}^* = S^*(\mathbf{e}_z \otimes \mathbf{e}_z - \frac{1}{3}\mathbf{I}) + T^*(\mathbf{e}_x \otimes \mathbf{e}_x - \mathbf{e}_y \otimes \mathbf{e}_y)$$

$$\mathbf{B}^* = S'^*(\mathbf{e}_z \otimes \mathbf{e}_z - \frac{1}{3}\mathbf{I}) + T'^*(\mathbf{e}_x \otimes \mathbf{e}_x - \mathbf{e}_y \otimes \mathbf{e}_y)$$

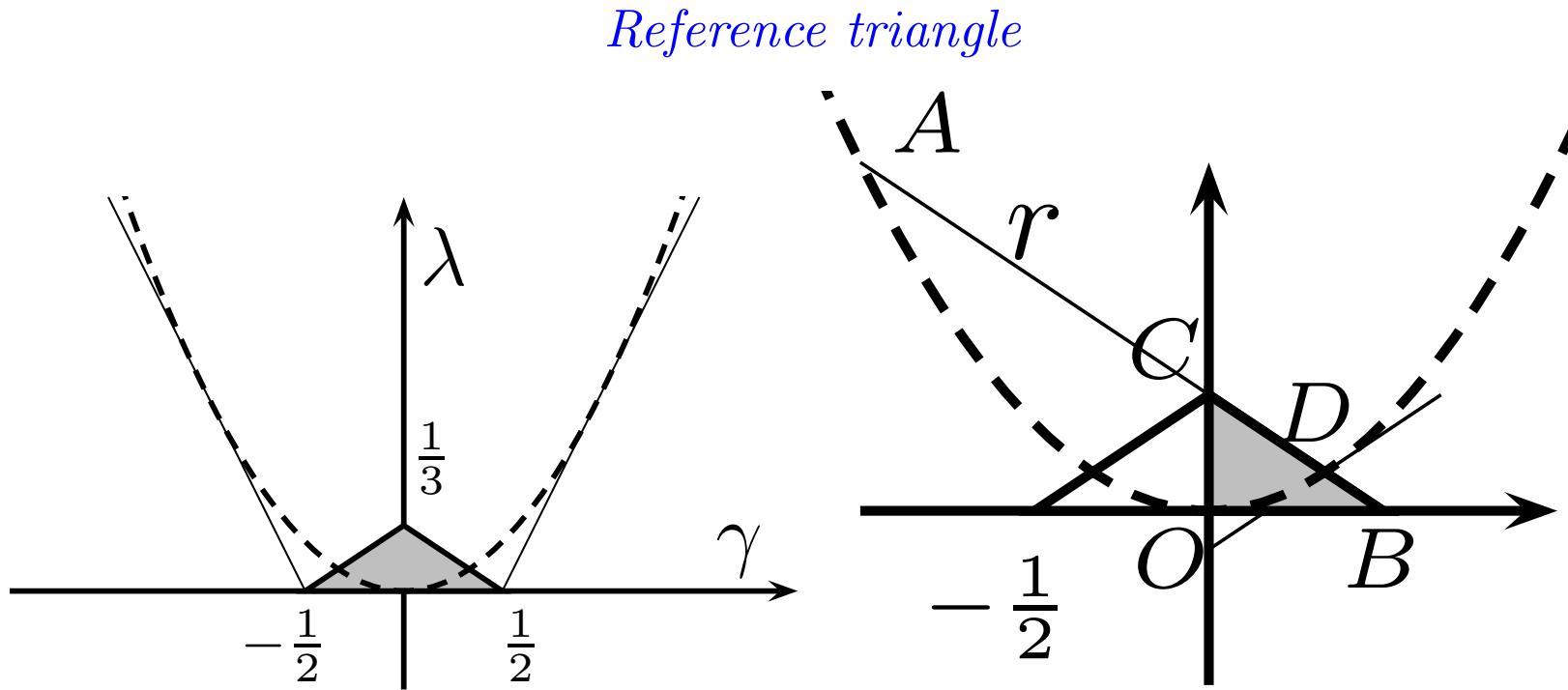
with the following correspondence on the order parameters

$$S = \frac{S'^* - S^*}{2}, \quad T = \frac{T'^* - T^*}{2},$$

$$S' = \frac{S'^* + 3S^*}{2}, \quad T' = \frac{T'^* + 3T^*}{2}$$

Conjugated phase diagram

$$2\gamma = 1 - 3\lambda \quad \frac{1}{9} \leq \lambda \leq 1$$



$$D\left(\frac{1}{3}, \frac{1}{9}\right), \quad B\left(\frac{1}{2}, 0\right), \quad C\left(0, \frac{1}{3}\right), \quad A(-1, 1)$$

$$r : 2\gamma + 3\lambda - 1 = 0$$

$2\gamma = 1 - 3\lambda, \quad 0 \leq \lambda \leq \frac{1}{9}$: a mildly repulsive model

$$V_\mu = -U_0 \left(\frac{1-\lambda}{2} \right) \{ \mathbf{b}^* \cdot \mathbf{b}^{*\prime} - \mu \mathbf{q}^* \cdot \mathbf{q}^{*\prime} \}$$

$$\mu = \frac{1-9\lambda}{1-\lambda}, \quad \beta_\mu = \beta \left(\frac{1-\lambda}{2} \right)$$

$$\lambda \in \left[0, \frac{1}{9} \right] \Leftrightarrow \mu \in [0, 1] \Leftrightarrow \overline{BD}$$

mean-field

$$f_\mu = \frac{1}{Z_\mu} \exp (\beta_\mu (\mathbf{B}^* \cdot \mathbf{b}^* - \mu \mathbf{Q}^* \cdot \mathbf{q}^*))$$

$$Z_\mu := \int_{\mathbb{T}} \exp (\beta_\mu (\mathbf{B}^* \cdot \mathbf{b}^* - \mu \mathbf{Q}^* \cdot \mathbf{q}^*))$$

$$\mathcal{F}_\mu (\mathbf{Q}^*, \mathbf{B}^*) = U_0 \left\{ \frac{1}{2} (\mathbf{B}^* \cdot \mathbf{B}^* - \mu \mathbf{Q}^* \cdot \mathbf{Q}^*) - \frac{1}{\beta_\mu} \log \left(\frac{Z_\mu}{8\pi^2} \right) \right\}$$

$$\mathbf{Q}^* = \langle \mathbf{q}^* \rangle = \int_{\mathbb{T}} f_\mu \mathbf{q}^* \quad \Leftrightarrow \quad \frac{\partial \mathcal{F}_\mu}{\partial \mathbf{Q}^*} = 0$$

$$\mathbf{B}^* = \langle \mathbf{b}^* \rangle = \int_{\mathbb{T}} f_\mu \mathbf{b}^* \quad \Leftrightarrow \quad \frac{\partial \mathcal{F}_\mu}{\partial \mathbf{B}^*} = 0$$

Bogoliubov's *minimax principle*

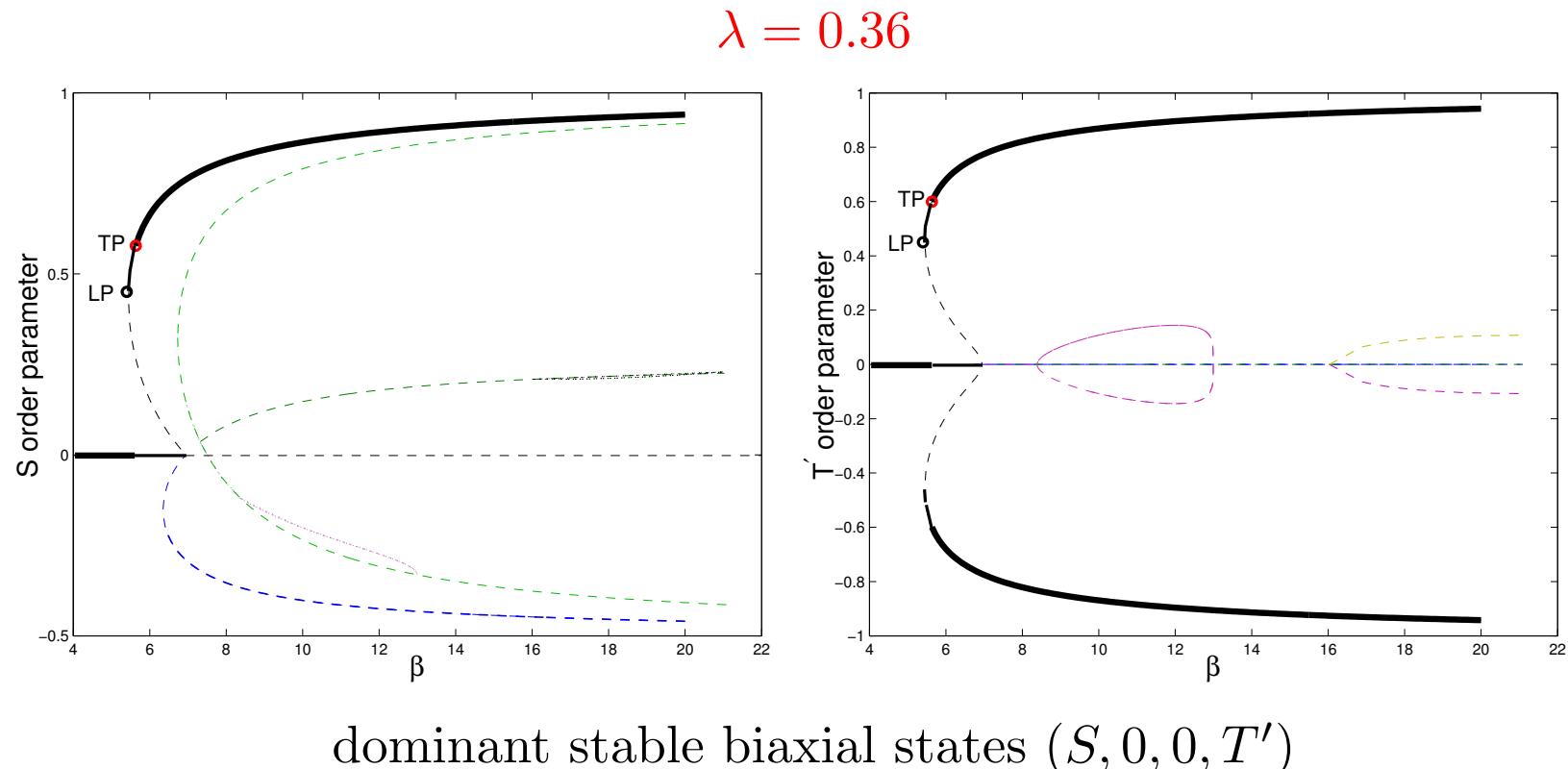
All the phases are characterized at least by two negative eigenvalues in the Hessian matrix of \mathcal{F}_μ . In this case we adopt a minimax principle to determine the effective phases.

$$\min_{\mathbf{B}^*} \max_{\mathbf{Q}^*} \mathcal{F}_\mu (\mathbf{Q}^*, \mathbf{B}^*)$$

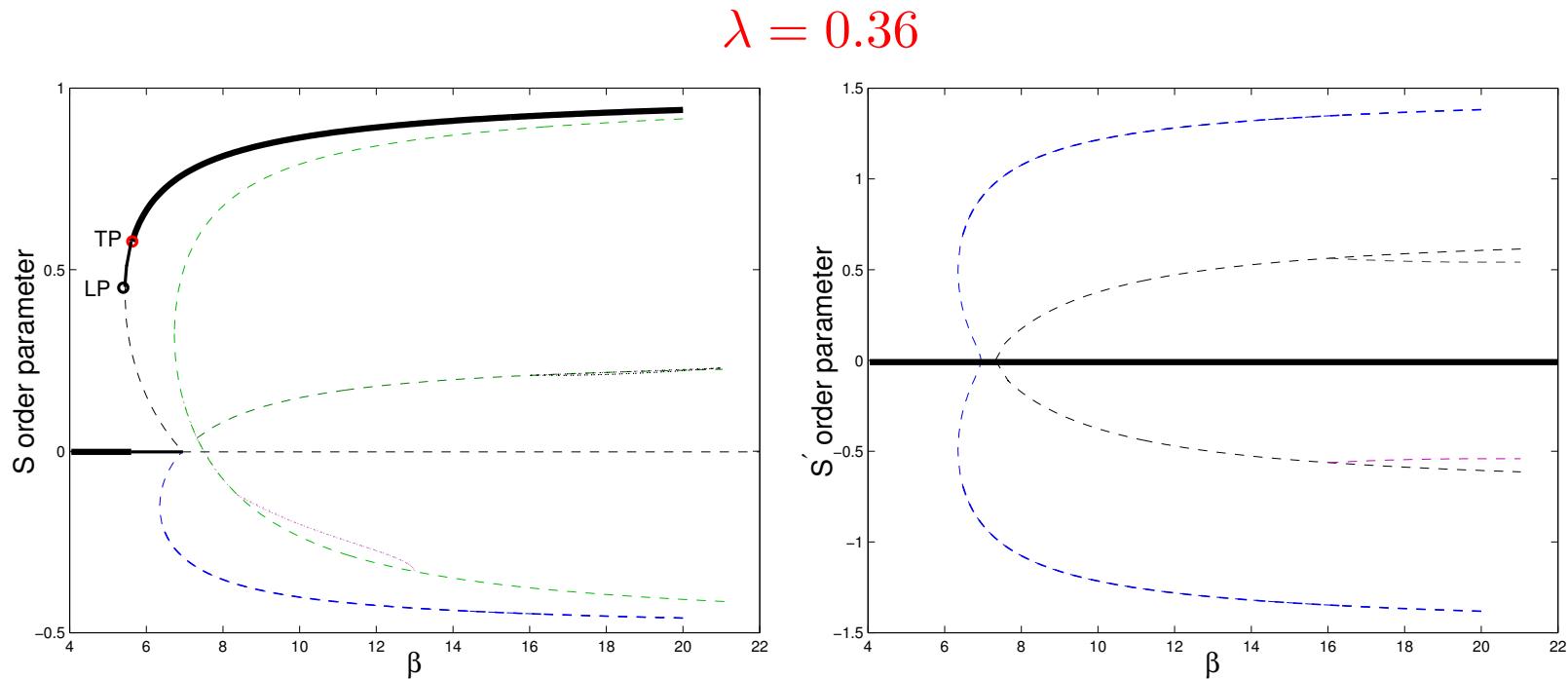
which is equivalent to finding a solution of the mean-field equations giving \mathcal{F}_μ a smaller value than do any other solutions

Future work

- Numerical continuation and Bifurcation Analysis for the mildly repulsive model (also with a γ -crossing term in the interaction)
- Landau theory for both attractive and mildly repulsive cases
.... and finally I would thank
- my PhD thesis advisor Prof. E. Virga
- Prof. Gartland, Prof. Durand, Prof. S. Romano and M. Osipov during my PhD research activity in Pavia



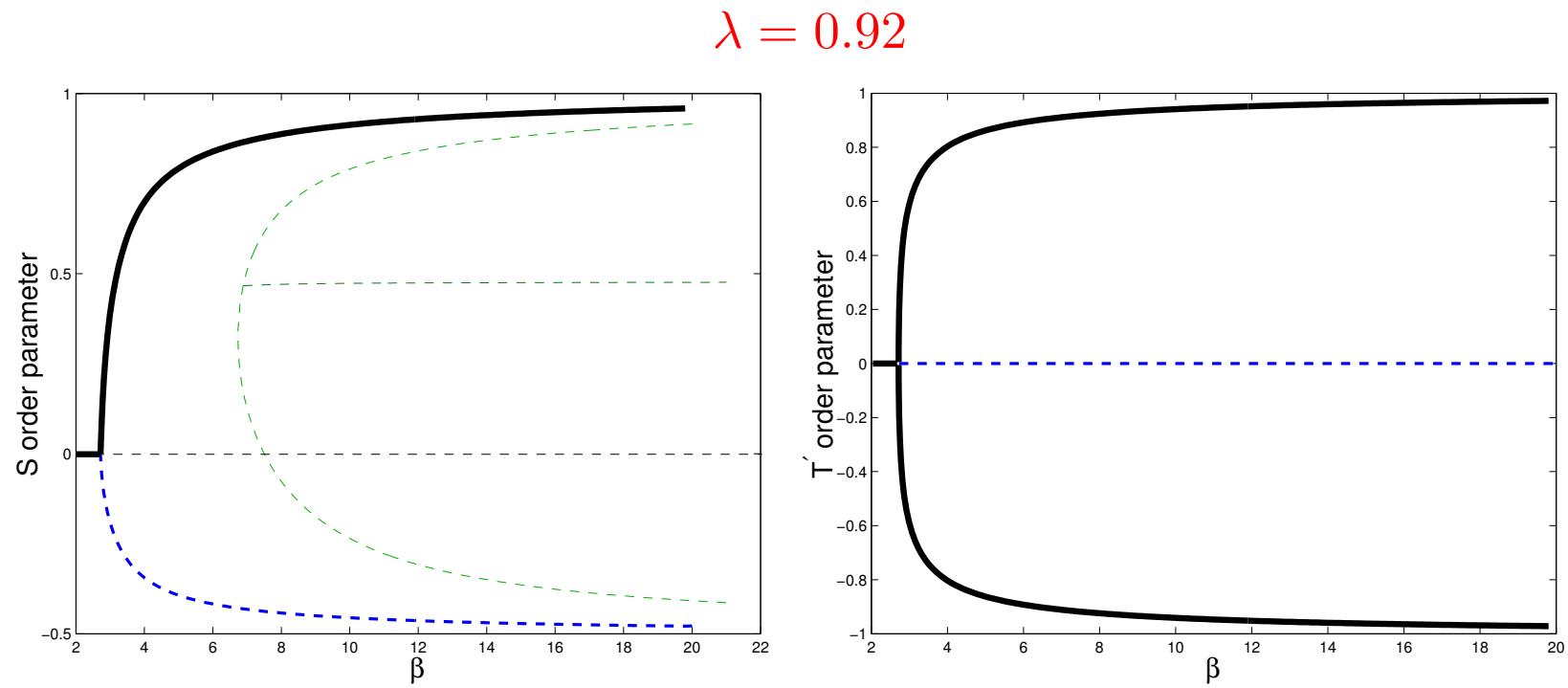
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unstable uniaxial states with a greater value of free energy $(S, 0, S', 0)$

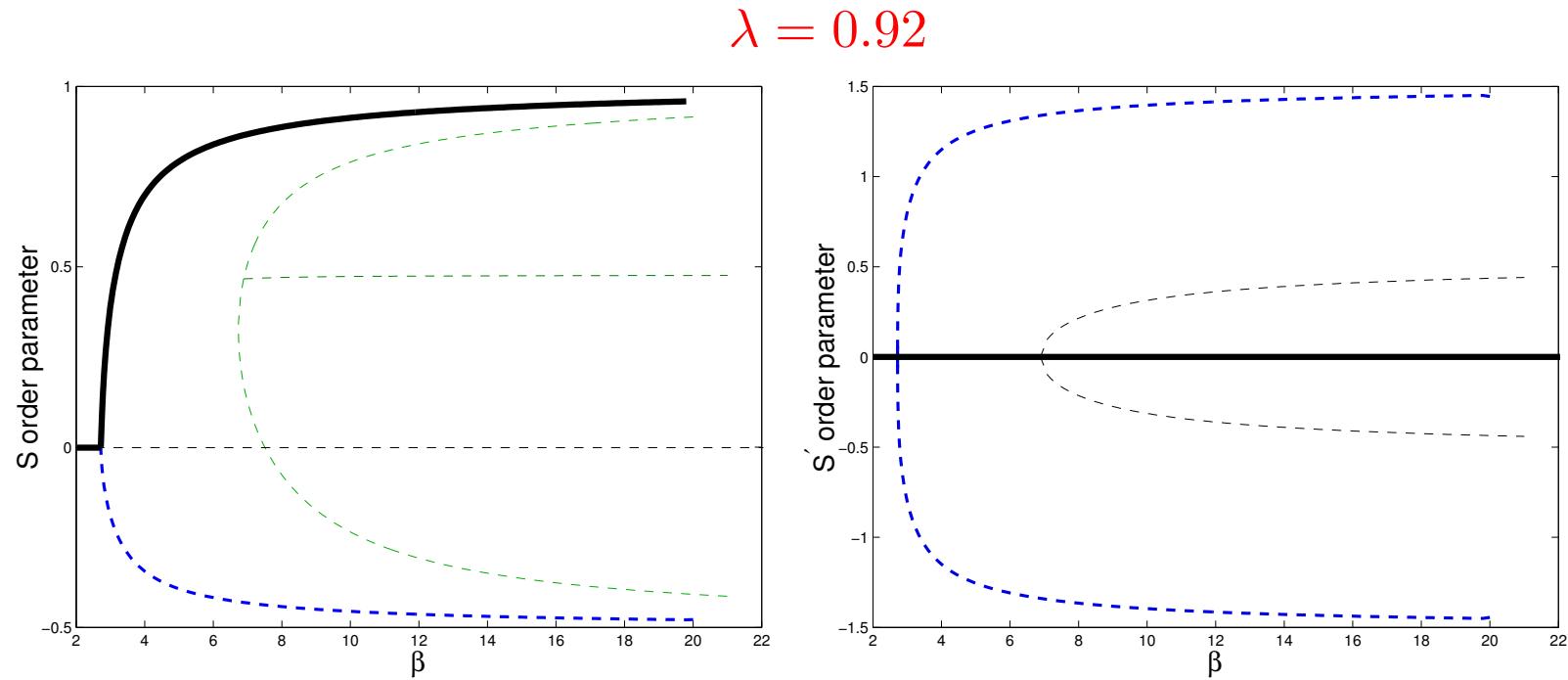
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dominant stable biaxial states $(S, 0, 0, T')$

$$\mathbf{Q} = S(\mathbf{e}_z \otimes \mathbf{e}_z - \frac{1}{3}\mathbf{I}) \quad \mathbf{B} = T'(\mathbf{e}_x \otimes \mathbf{e}_x - \mathbf{e}_y \otimes \mathbf{e}_y)$$



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